

LINE PARAMETERS OF VOLCANIC GASES (CO₂ AND CO): MEASUREMENTS AND CALCULATIONS



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Introduction

When studying volcanic activity, knowledge of the carbon oxides dynamics in the surface layers of the atmosphere is necessary. These studies are conducted when recording solar atmospheric spectra using ground stations equipped with FTS spectrometers. To record the solar spectra, we need data on the broadening coefficients of the lines of CO, CO₂, and H₂O broadened by the main atmospheric gases, such as nitrogen, oxygen, and carbon dioxide, the latter being released during volcanic activity.

The absorption of ¹²C¹⁶O₂ in the spectral region 9340–9650 cm⁻¹ was measured on a Bruker IFS 125HR Fourier spectrometer, combined with a multi-pass measuring cell. A large array of the coefficients of broadening and shift of the lines of the 20032–00001 band was obtained.

The broadening coefficients of the carbon monoxide and dioxide lines, as well as the water vapor lines were calculated. Two calculation methods were used: a semi-empirical method, based on the semiclassical impact theory of line broadening, and supplemented by the introduction of a correction factor, the parameters of which were determined using experimental data [Mol. Phys. 102 (2004)], as well as the average frequencies method [Opt. Atmosphere Ocean. 28 (2015)]. The lines half-widths were calculated for room temperature (296 K) and for the Earth's atmosphere temperature range. The calculations were performed for a wide range of rotational quantum numbers J (up to 100). For each half-width the temperature exponents were calculated. Experimental and calculated are in a good agreement.

Semi-empirical approach

The semi-empirical method [1] simplifies the semi-classical Robert-Bonamy expressions to an Anderson-type form via introducing effective correction factors. These factors, chosen in a convenient parameterized form, account for the real trajectory curvature, vibrational dependence and corrections to the scattering matrix. The model parameters are fitted on some experimental data and their values are further used to compute the line-widths for other sets of rotational quantum numbers (the vibrational dependence of line-broadening coefficients is negligible for most molecular systems and the parameters adjusted on one vibrational or rotational band can be safely used for other vibrational bands).

The semi-classical half-width γ_{if} associated with the optically active $i \rightarrow f$ transition are given by

$$\gamma_{if} - i\delta_{if} = \frac{n_2}{c} \sum_{J_2} \rho(J_2) \int_0^\infty d\nu \nu f(\nu) \int_0^\infty db b S(b)$$

with the interruption function

$$S(b) = 1 - \sum_{J_2, m, m'} \frac{(J_2, m, q, J_2, m_1)(J_2, m, q, J_2, m_1)}{(2J_2 + 1)(2J_2 + 1)} (J_2, m, J_2, m_2, T | J_2, m_1, J_2, m_2) \times \\ \times (J_2, m', J_2, m_2, T | J_2, m', J_2, m_2)$$

The half-width can be expressed as

$$\gamma_{if} = A(if) + \sum_{i'l} D^2(ii'l)P_i(\omega_{i'l}) + \sum_{f'l} D^2(ff'l)P_f(\omega_{f'l}) + \dots$$

where
$$A(if) = \frac{n_2}{c} \sum_{J_2} \rho(J_2) \int_0^\infty \nu F(\nu) b_0^2(\nu, J_2, i, f) d\nu$$

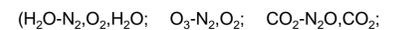
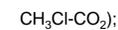
The efficiency function $P(\omega)$ can be represented as

$$P_i(\omega) = \frac{n_2}{c} \sum_k \rho(J_2) \sum_{l, J_2'} A_{il} D^2(J_2, J_2' | l) F_{il} \left(\frac{2\pi c b_0(J_2, i, f)}{\nu} (\omega + \omega_{J_2, J_2'}) \right)$$

$$P_i(\omega_{i'l}) = P_i^l(\omega_{i'l}) C_l(\omega_{i'l})$$

where correction factor $C_l(\omega)$ is usually taken in the following forms:

$$1) \quad C_l(\omega_{i'l}) = \frac{c_1}{c_2 \sqrt{J+1}} \quad \text{for linear molecules, symmetric and asymmetric tops}$$



Temperature dependence

The temperature-dependence characteristics were determined for the interval 200–400 K.

The usual relation for temperature dependence of broadening and shifting coefficients are given by

$$\gamma(T) = \gamma(T_0) \left(\frac{T_0}{T} \right)^N$$

Basic idea

Upper and lower state quantum numbers

Averaged energy differences of coupled states

Line broadening coefficients

Averaged energy difference approach

Based on the modified RB formalism, the expressions of the half-widths are given by

$$\gamma = \frac{n_2}{2\pi c} \int_0^\infty \nu f(\nu) d\nu \int_0^\infty 2\pi b db [1 - \cos(S_1 + \text{Im} S_2)] e^{-\text{Re} S_2} \\ \approx \frac{n_2 \nu}{2\pi c} \int_0^\infty 2\pi b db (1 - e^{-\text{Re} S_2}), \quad S_{2, \text{outer}, i} = \sum_{L_1, K_1, K_1'} \int_0^t dt dt' G_{L_1, K_1, K_1'}(t, t') W_{L_1, K_1, K_1'}^{(a)}(t - t') W_{L_2}^{(b)}(t - t')$$

In the above expression, two functions which are independent of the potential are defined by

$$W_{L_1, K_1, K_1'}^{(a)}(t) = \sum_{i, i'} (2i' + 1) D(i\tau i'\tau'; L_1, K) D(i\tau i'\tau'; L_1, K') e^{i\omega_{i\tau i'\tau'} t}$$

where quantum indices i and i' associated with the H₂O molecule represent pairs of the quantum numbers J_i, τ_i and $J_{i'}, \tau_{i'}$, used to specify the energy levels of the initial and final states. Where $D(i\tau i'\tau'; L, K)$ is defined by

$$W_{L_2}^{(b)}(t) = \sum_{i_2, i_2'} (2i_2 + 1)(2i_2' + 1) \rho_{i_2} C^2(i_2 i_2' L_2, 000) e^{i\omega_{i_2 i_2'} t}$$

The functions introduced above are:

$$G_{L_1, K_1, K_1'}(t, t') = \frac{\hbar^{-2}}{4\pi(2L_1 + 1)^2(2L_2 + 1)^2} (-1)^{K_1 + K_1'} \sum_{L_2} (-1)^{L_2 + L_2'} (2L_2 + 1) \\ \times u(L_1, L_2, L_2'; K_1; R(t)) u(L_1, L_2, L_2'; K_1'; R(t')) P_L(\cos \Theta_{i, i'})$$

$$\text{Re}(S_{2, \text{outer}, i}) = \sqrt{\frac{\pi}{2}} \sum_{L_1, K_1, K_1'} \sum_{L_2, i_2} \{ \sum_{i_2'} (2i_2 + 1)(2i_2' + 1) \rho_{i_2} C^2(i_2 i_2' L_2, 000) \} \\ \times \{ \sum_{i_2'} (2i_2' + 1) D(i\tau i'\tau'; L_1, K_1) D(i\tau i'\tau'; L_1, K_1') \} H_{L_1, K_1, K_1'}(\omega_{i\tau i'\tau'} + \omega_{i_2 i_2'})$$

Then, we can estimate how values of $\text{Re}(S_2)$ vary with different lines of interest. We introduce coupled state energy differences between a H₂O line of interest and other coupled lines defined

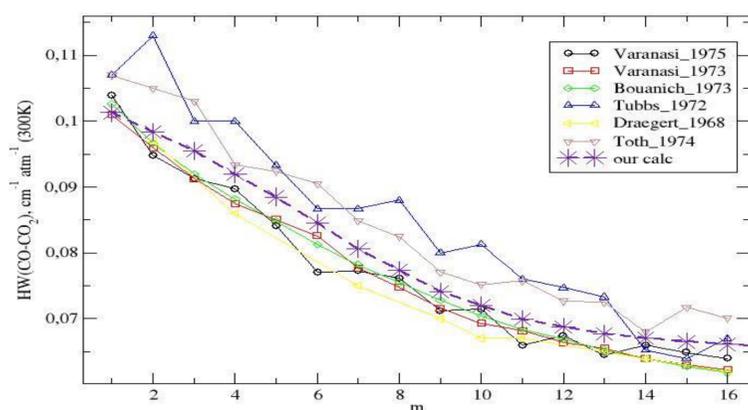
$$\bar{\omega}_{if} = \frac{1}{W} \left\{ \sum_{i'} (2i' + 1) D^2(ii'l) f(\omega_{i'l}) \omega_{i'l} + \sum_{f'} (2f' + 1) D^2(ff'l) f(\omega_{f'l}) \omega_{f'l} \right\} \equiv \omega'$$

by
$$W = \sum_{i'} (2i' + 1) D^2(ii'l) f(\omega_{i'l}) + \sum_{f'} (2f' + 1) D^2(ff'l) f(\omega_{f'l})$$

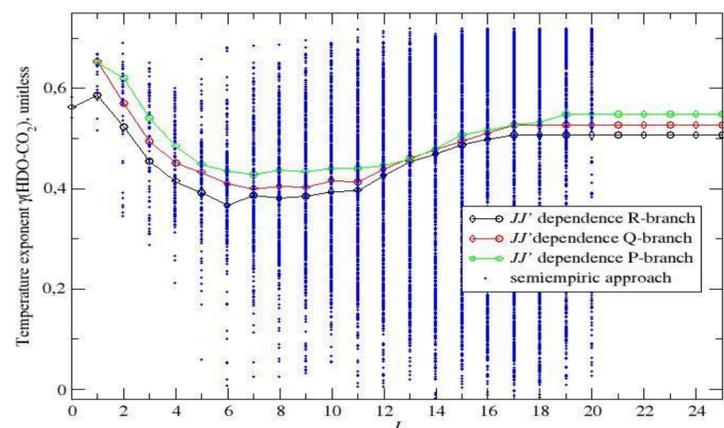
Calculations were made considering contributions of the dipole-quadrupole and quadrupole-quadrupole interactions.

LINE BROADENING: CO-CO₂, CO₂-CO₂, H₂O-CO₂

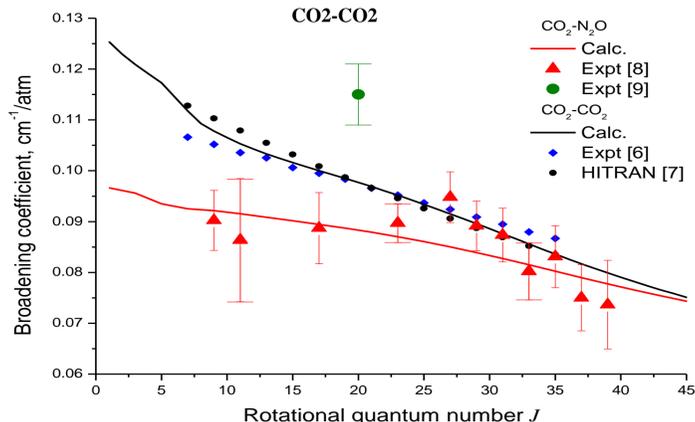
CO-CO₂



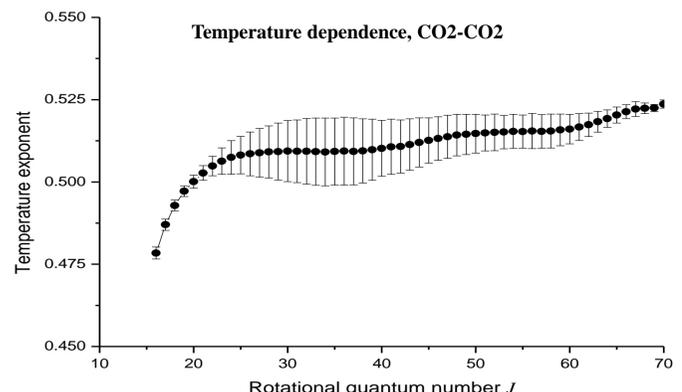
H₂O-CO₂



CO₂-CO₂



Temperature dependence, CO₂-CO₂



References

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Conclusions

Measurements and calculations of CO₂ broadening coefficients and their temperature exponents for CO₂, CO and H₂O lines were obtained for the temperature interval 150–350 K recommended for HITRAN. Calculations are performed using two methods: 1. the average energy difference method, 2. the semi-empirical approach, which is a modification of the impact theory performed by introduction of the empirical correction factor. Data are obtained for a large amount of rovibrational transitions J and K_a requested by spectroscopic databases. Theoretical coefficients agree satisfactorily with measured values. The calculations have demonstrated that the semi-empirical method provide a reasonable description of the variation of the broadening coefficients and their temperature exponents.

The capacity of the semi-empirical approach to reproduce available measurements of line-shape parameters and to extrapolate their values to higher J and K_a inaccessible experimentally can be also exploited to get line-shape characteristics for another spectral range of considered colliding system, such calculations for the other bands are in progress. We have computed line-lists for broadening coefficients as well as the corresponding temperature exponents for the temperature interval 200–400 K recommended for HITRAN.

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